# Particle Acceleration and Emission Regions of Gamma-Ray Pulsars

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#### Outline

- 1. Introduction
- 2. Acceleration region : Pair-starved MSPs
- 3. Emission region : sub-TeV emission region

### Structure of pulsar magnetosphere

Rotating magnet  $\rightarrow$  Unipolar induction  $\rightarrow E_{||}$  screening  $\rightarrow$  Magnetosphere

**Basic concept** (Goldreich & Julian 1969) Θ Light cylinder Θ Θ Θ Null line Θ Θ Θ Ð Ŧ Ð Ð

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- Particle supply (particle creation)
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  - $\rightarrow$  GeV  $\gamma$ -ray (e<sup>±</sup> CR ?)
- Particle supply (particle creation)
  - $\rightarrow$  IR ~ X-ray (& TeV  $\gamma$ -ray ?)
- E<sub>||</sub> screening (ρ & J adjustment)
   → Radio ?

#### Problems

 Outer gap model cannot reproduce observed γ-ray light curves in some millisecond pulsars (MSPs).

 VERITAS and MAGIC detected sub-TeV pulsed emission from Crab pulsar.

#### Unusual γ-ray light curve

Two MSPs show unusual light curves that the γ-ray peak lead the radio peak. Young γ-ray pulsars do not show such light curves.



Guillemot (2009)

Pair-Starved Polar Cap (1/2) Muslimov & Harding (2004) suggest that most MSPs can not produce pairs through curvature radiation. In this case, particle acceleration can occur in almost open field volume (Pair-Starved Polar Cap model).



#### Pair-Starved Polar Cap (2/2)

The model that have extended particle acceleration region can explain observed light curves.



Conditions for radio-emitting region In general, people have believed the conditions of a primary beam with  $\gamma_b \sim 10^7$  and  $n_b \sim n_{GJ}$ , and a secondary  $e^{\pm}$  with  $\gamma_p \sim 10-10^3$  and  $n_p \sim 10^3-10^5 n_{GJ}$ in the radio-emitting region.



However, the existence of pair-starved MSPs suggests that radio emission mechanisms should be insensitive to the particle number density down to sub-GJ one. Therefore, further verification is important.

#### Acceleration

#### Assumptions

- In the wind region,
- the energy equipartition between B and e<sup>±</sup>.
- the flux conservation of B and e<sup>±</sup>.

Typical energy is related to the number of particles.

$$\varepsilon_e = e\Delta V_{\max}\kappa^{-1} \sim 50\kappa^{-1}$$
TeV

Propagation of e<sup>±</sup> cosmic ray  $\frac{\partial}{\partial t}f(t, r, \epsilon_e) = D(\epsilon_e)\nabla^2 f + \frac{\partial}{\partial \epsilon_e}\left(P(\epsilon_e)f\right) + Q(t, r, \epsilon_e)$ Diffusion coefficient  $D(\varepsilon_e) = D_0 (1 + \varepsilon_e/3 \text{GeV})^{\delta}$ **Cooling function**  $P(\varepsilon_e) = \frac{4\sigma_T \varepsilon_e^2}{3m_e^2 c^3} \left[ \frac{B^2}{8\pi} + \int d\varepsilon_\gamma u_{\text{tot}}(\varepsilon_\gamma) f_{\text{KN}}\left( \frac{4\varepsilon_e \varepsilon_\gamma}{m_e^2 c^4} \right) \right]$ Source function  $Q_0(\varepsilon_e, \tilde{t}) \propto \varepsilon_e^{-\alpha} \exp\left(-\frac{\varepsilon_e}{\varepsilon_{\text{cut}}}\right) \left(1 + \frac{\tilde{t} - t_i}{\tau}\right)^{-2}$ Assumptions MSPs are single population. Source function is mono-energetic distribution. Spectrum 

$$f_{\text{ave}}(\varepsilon_e) = \int_0^{t_0} dt_i \int_0^{a_{\text{diff}}(\varepsilon_e, \varepsilon_{e,i})} 2\pi r dr f(t_0, r, \varepsilon_e; t_i) R$$
  
R ~ 3×10<sup>-9</sup> kpc<sup>-2</sup>yr<sup>-1</sup> : local birth rate



![](_page_13_Figure_0.jpeg)

Do Sub-TeV emission regions locate inside the light cylinder ? or outside ?

#### γ-B absorption Lee et al. (2010)

![](_page_14_Figure_1.jpeg)

The value of the last escaping radius for a 400(120) GeV photon is  $0.5(0.35) R_{lc}$ .

#### Geometrical modeling Assumptions

- Magnetic field : Rotating dipole
- Emission direction : Along the particle trajectory
- Emission region : r<sub>null</sub> < r < R<sub>lc</sub> (Outward)

Emissivity

: Constant

:α

: 7

 $r_{null}$ : Distance to null surface  $R_{lc}$  : Light cylinder radius

f=1.0

We focus on the peak phases of the light curve.

Parameters

- Magnetic colatitude :  $f \Xi r_{pc}/r_{pc,0}$
- Inclination angle
- Viewing angle

#### Emission region in outer gap

The distance distribution

from NS to emission regions

![](_page_16_Figure_3.jpeg)

Peaks include photons emitted from relatively far region.

#### Emission region in outer gap

The distance distribution

from NS to emission regions

![](_page_17_Figure_3.jpeg)

Φ

Peaks include photons emitted from relatively far region.

All photons at leading side of second peak come from the closest region in the magnetosphere.

![](_page_18_Figure_0.jpeg)

![](_page_19_Figure_0.jpeg)

We can exclude the photons from outer region when we use photons in phase < φ\_peak - 0.05.</li>
Photons from the closest region from NS are included in phase φ\_peak - (0.05 - 0.15).

#### Constraint for outer gap model

![](_page_20_Figure_1.jpeg)

#### Prospects

- 1. If the spectrum have the power-law component... outer gap  $\rightarrow$  Constraint for geometrical parameters.
- 2. If the spectrum show the strong cut-off... outer gap → Constraint for the altitude of emission region. cold wind → Constraint for the anisotropy of the particle distribution in pulsar wind.
- 3. If the spectrum do not have the power-law component... outer gap → Constraint for the condition of IC. cold wind → Constraint for the anisotropy of the particle distribution in pulsar wind.

#### Summary

- Outer gap model cannot reproduce observed γ-ray light curve in some MSPs.
- → CTA can detect the evidence of the existence of pair-starved MSPs in the e<sup>±</sup> cosmic ray spectrum.
- VERITAS and MAGIC detected sub-TeV emission from Crab pulsar.
- → CTA can restrict the emission region using the spectrum at the phase of the bridge.